

Formal Analysis of Quantum Systems using Process Calculus

Simon Gay University of Glasgow









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I will try to answer the following questions.

- What is quantum information processing?
- Why is it interesting?
- Why do I want to apply formal methods?
- What can be done with quantum process calculus?



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QFOX: Quantum Computation: Foundations, Security, Cryptography and Group Theory (EPSRC 2008-2011)

Lord Kelvin / Adam Smith PhD Scholarship (University of Glasgow)

QUISCO (Quantum Information Science Scotland)



The idea is to represent information by means of physical systems whose behaviour must be described by quantum physics, and to process information by means of operations that arise from quantum physics.

Examples: the spin state of a single atom the polarization state of a single photon



Superposition

The state of a classical bit is either 0 or 1. The state of a quantum bit (qubit) is

$$lpha |0
angle$$
 + $eta |1
angle$

where |0
angle and |1
angle are the basis states.

For example:
$$\frac{1}{\sqrt{2}}|0\rangle + \frac{1}{\sqrt{2}}|1\rangle$$
 or $\frac{\sqrt{3}}{2}|0\rangle - \frac{1}{2}|1\rangle$

A qubit may be in a basis state or it may be in a superposition state.



Measurement

It is not possible to inspect the components of a quantum state. We can only do a measurement.

Measuring a qubit that is in state $\alpha |0\rangle + \beta |1\rangle$

has a random result:

with probability $|lpha|^2$ the result is $|0\rangle$ with probability $|eta|^2$ the result is $|1\rangle$

Operations on a Superposition

An operation acts on every basis state in the superposition. For example, if we have the state

$$\frac{1}{2}|000\rangle + \frac{1}{2}|010\rangle - \frac{1}{2}|110\rangle - \frac{1}{2}|111\rangle$$

and we apply the operation "invert the second bit" then we get the state

$$\frac{1}{2}|010\rangle + \frac{1}{2}|000\rangle - \frac{1}{2}|100\rangle - \frac{1}{2}|101\rangle$$

Note that a measurement does **not** allow us to discover the result of applying the operation to every basis state.



No Cloning

It is not possible to define an operation that reliably makes a duplicate of an unknown quantum state.

(Contrast this with the classical situation, where the existence of uniform copying procedures is one of the main advantages of digital information.)

It is possible to transfer an unknown quantum state from one physical carrier to another, but the process destroys the original. This is known as quantum teleportation.



Entanglement

The states of two (or more) qubits can be correlated in a way that is stronger than any possible classical correlation. This is known as quantum entanglement.

Many quantum algorithms and communication protocols make use of entanglement.

$$\frac{1}{\sqrt{2}}|00\rangle + \frac{1}{\sqrt{2}}|11\rangle$$



Deutsch-Jozsa Algorithm (David Deutsch / Richard Jozsa, 1992)

Given a function $f: \{0,1\}^n \rightarrow \{0,1\}$ that is either constant or balanced, works out which, with only one evaluation of f.

A classical algorithm would require $2^{n-1} + 1$ evaluations in the worst case.









Shor's Algorithm (Peter Shor, 1994)

A quantum algorithm for factorising integers, which is more efficient than any known classical algorithm.

 $(\log n)^3$ vs. $e^{(\log n)^{1/3}(\log \log n)^{2/3}}$

The RSA cryptosystem used for internet security relies on the assumption that factorisation is difficult. A practical implementation of Shor's algorithm would threaten current information security technology

However, factorisation is believed not to be a member of the class of intractable problems known as NP-complete problems.







Grover's Algorithm (Lov Grover, 1996)

A quantum search algorithm. It requires \sqrt{n} steps to find an item in an unstructured list of length n.

Classically, every item must be inspected, requiring n/2 operations on average.





BB84 (Charles Bennett & Gilles Brassard, 1984)

A protocol for generating shared cryptographic keys. Its security relies only on the laws of quantum physics, especially the no-cloning principle. It is secure against all possible future developments in quantum computing.







It is interesting to understand the information-processing power permitted by the laws of physics.

Quantum algorithms might help to solve some classes of problem more efficiently.

But if NP-complete problems cannot be solved efficiently even by a quantum computer, understanding why not is also of fundamental interest.

Quantum cryptography deals with any threat that quantum computing poses to classical cryptography.



As integrated circuit components become smaller, quantum effects become difficult to avoid. Quantum computing might be necessary in order to continue making computers smaller.

Richard Feynman (1982) suggested that quantum computers could be used to simulate complex (quantum) physical systems, whose behaviour is difficult to analyse.













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Most researchers are not convinced that D-Wave has built a quantum computer.

D-Wave is pioneering the development of a new class of high-performance computing system designed to solve complex search and optimization problems, with an initial emphasis on synthetic intelligence and machine learning applications.

D-Wave systems are architected around an innovative processor that uses a computational model known as adiabatic quantum computing (AQC). These processors exploit quantum effects to solve search and optimization problems in a new way. They are fabricated using superconducting metals instead of semiconductors and are operated at ultra-low temperatures in a magnetic vacuum.

AQUA@HOME

D-Wave's AQUA@home (Adiabatic QUantum Algorithms) is a research project whose goal is to predict the performance of superconducting adiabatic quantum computers on a variety of hard problems arising in fields ranging from materials science to machine on parning. AQUA@home uses Internet-connected IS computers to help design and a analyze quantum computing algorithms, using Quantum Monte Carlo techniques. You can participate by running a free program on your computer.

AQUA@home white paper AQUA@home website



Quantum cryptography is already practical. Whether there is real demand for it remains to be seen.

Quantum computing seems feasible in principle, although there are formidable scientific and engineering problems. Decoherence: loss of quantum state due to unwanted interaction with the environment.

BUT remember that in 1949, the statement "In the future, computers may weigh no more than 1.5 tonnes" was a speculative prediction.



There is no doubt about the correctness of quantum algorithms and protocols.

Teleportation can be checked with a few lines of algebra.

Shor's and Grover's algorithms have been thoroughly studied.

Mayers (2001) and others have proved the security of quantum cryptography.

But what about systems – combinations of classical and quantum communication and computation?



Raja Nagarajan and I (2002) suggested applying formal methods to quantum systems, with the same motivation as for classical systems:

Formal modelling languages, for unambiguous definitions.

Analysis of systems, rather than idealized situations.

Systematic verification methodology, rather than ad hoc reasoning.

Tool support.



I have been working on the following strands:

Process calculus for quantum systems (with Raja Nagarajan and Tim Davidson)

Model-checking for quantum systems (with Raja Nagarajan and Nick Papanikolaou)

Formal description of physics experiments (with Sonja Franke-Arnold and Ittoop Vergheese Puthoor)

For the rest of this talk, I will focus on process calculus.



Quantum Teleportation

A protocol for transferring an unknown quantum state from Alice to Bob, making use of classical communication and some pre-existing shared entanglement.





We have developed the process calculus **Communicating Quantum Processes** (CQP) for modelling combined quantum/classical systems (SG + Nagarajan, 2005, 2006).

CQP has a formal syntax, a formal operational semantics, and a type system.

(Other approaches: QPAIg (Jorrand+Lalire 2004), qCCS (Ying et al. 2006-2011))



Quantum Teleportation in CQP



Alice(q,in,out) = in?[u] . { u,q *= CNot } . { u *= H } . out![measure u,q] . Stop





A configuration consists of a quantum state and a process.

The semantics of CQP specifies transitions between configurations.

Each transition is a communication or an operation on the quantum state.

 $w = \alpha |0\rangle + \beta |1\rangle$; a![w]. Stop | Teleport(a,b) | b?[v]. Stop

=

w = α|0>+β|1> ; a![w] . Stop | (qbit x, y)({ x *= H } . { x,y *= CNot } . (new c)(Alice(x,a,c) | Bob(y,c,b))) | b?[v] . Stop



w,x,y = $(\alpha|0\rangle+\beta|1\rangle)|00\rangle$;

a![w] . Stop | { x *= H } . { x,y *= CNot } . (new c)(Alice(x,a,c) | Bob(y,c,b)) | b?[v] . Stop

V

=

```
w,x,y = (\alpha|0\rangle+\beta|1\rangle)(|00\rangle+|11\rangle);
```

a![w] . Stop | (new c)(Alice(x,a,c) | Bob(y,c,b)) | b?[v] . Stop

w,x,y = (α|0>+β|1>)(|00>+|11>) ; a![w] . Stop | a?[u] . { u,x *= CNot } . { u *= H } . c![measure u,x] . Stop | Bob(y,c,b) | b?[v] . Stop



w,x,y = $(\alpha|0\rangle+\beta|1\rangle)(|00\rangle+|11\rangle)$;

{ w,x *= CNot } . { w *= H } . c![measure w,x] . Stop | Bob(y,c,b) | b?[v] . Stop

$\mathbf{\Psi}$

 $w,x,y = \alpha(|000\rangle + |011\rangle + |100\rangle + |111\rangle) + \beta(|001\rangle + |010\rangle - |101\rangle - |110\rangle) ;$

c![measure w,x] . Stop | Bob(y,c,b) | b?[v] . Stop

$\mathbf{\Psi}$

▶ 1⁄4



w,x,y = $\alpha|011\rangle+\beta|010\rangle$; c![1]. Stop | Bob(y,c,b) | b?[v]. Stop

w,x,y = α |011>+ β |010> ; c![1] . Stop | c?[r] . { y *= Pauli_r } . b![y] . Stop | b?[v] . Stop L $w,x,y = \alpha |011\rangle + \beta |010\rangle$; { y *= Pauli₁ } . b![y] . Stop | b?[v] . Stop J w,x,y = $\alpha |010\rangle + \beta |011\rangle$; b![y]. Stop | b?[v]. Stop = w,x,y = $|01\rangle(\alpha|0\rangle+\beta|1\rangle)$; b![y]. Stop | b?[v]. Stop



We defined Teleport(a,b), which receives a qubit on channel a and sends a qubit on channel b, using teleportation in between.

The following process has the same effect, and we regard it as a **specification** of teleportation.

Identity(a,b) = a?[x] . b![x] . Stop

Now we want to state the requirement that

 $Teleport(a,b) \approx Identity(a,b)$

and prove that it is satisfied. So we need a theory that defines ≈ and provides some proof techniques.



The relation \approx is a behavioural equivalence: if $P \approx Q$ then P and Q have indistinguishable behaviour.

It is a form of probabilistic branching bisimulation, where the observations take into account the amount of information that a transition reveals about the quantum state.

(Matching of transitions considers the reduced density matrix w.r.t. input/output qubits).

The aspiration for behavioural equivalence is **congruence**:

 $P \approx Q \implies C[P] \approx C[Q]$

for all process contexts C. This supports equational reasoning. Congruence properties are sometimes known as composability properties.



Obtaining a congruence relation for a quantum process calculus was an open problem for a while.

We have solved it for CQP (Tim Davidson's PhD thesis, 2011) and Ying et al. have independently solved it for qCCS (POPL 2011). (The details are quite complicated and require changes to the semantics).

We can show that $Teleport(a,b) \approx Identity(a,b)$ and therefore this equivalence holds in all contexts.

Although correctness of teleportation is standard, this formulation is (we claim!) a valuable new perspective.



A mixed quantum state is a probability distribution over pure quantum states, representing classical uncertainty.

We introduce mixed CQP configurations, distinct from probabilistic configurations. A measurement produces a mixed configuration. If the measurement result is output then a probabilistic configuration is produced.

Internal communication of a measurement result, however, does not remove mixedness.

In teleportation, with this new semantics, there are no probabilistic configurations, because the measurement result is never output.



Consider $P = a?[x] . \{measure x\} . Stop$ $Q = a?[x] . {x *= H} . {measure x}. Stop$

They are equivalent, in all quantum states, just because they produce no output.

Put them in parallel with R = b![q]. Stop in the state $p,q = |00\rangle + |11\rangle$.

If the measurement produces a probabilistic configuration, and R outputs afterwards, then the possible reduced density matrices for q, produced by $P \mid R$ and $Q \mid R$, are different.

This means that $P \mid R$ and $Q \mid R$ do not have matching output transitions.



Consider $P = a?[x] . \{measure x\} . Stop$ $Q = a?[x] . \{x *= H\} . \{measure x\}. Stop$

They are equivalent, in all quantum states, just because they produce no output.

Put them in parallel with R = b![q]. Stop in the state $p,q = |00\rangle + |11\rangle$.

In the modified semantics, the measurement produces a mixed configuration and because the result is not output, it never becomes a probabilistic configuration.

Then the output of q has the same reduced density matrix for both $P \mid R$ and $Q \mid R$.



To make reasoning about processes easier and more practical:

equational axiomatization of equivalence

automated equivalence checking

More substantial applications, e.g. cryptography.



Quantum information processing is a fascinating research area which might lead to important computational and cryptographic technologies.

In any case, seeking to understand the computational power of quantum systems is a basic research question that approaches fundamental physics from an interesting new angle.

The formal methods approach, and process calculus in particular, will be needed for assurance of practical systems, and gives an interesting new perspective on quantum behaviour.